

Reminder of our last lecture

Tangent vectors

A tangent vector $V_p \in T_p M$ is a map $V_p: C^\infty(M; \mathbb{R}) \rightarrow \mathbb{R}$ that fulfills for all $a, b \in \mathbb{R}$ and $f, g \in C^\infty(M; \mathbb{R})$ that:

(i. Linearity)

$$V_p[af + bg] = aV_p[f] + bV_p[g]$$

(ii. Leibniz rule)

$$V_p[f \cdot g] = V_p[f]g(p) + f(p)V_p[g]$$

We also saw:

(i.) $T_p M$ is a real vector space of dimension $\dim(T_p M) = \dim(M) = n$

(ii.) For any smooth path $\gamma: [-1, 1] \rightarrow M; t \mapsto \gamma(t)$ passing through $p = \gamma(0)$ the equation

$$\dot{\gamma}_p[f] = \left. \frac{d}{dt} \right|_{t=0} (f \circ \gamma)(t)$$

defines a tangent vector

(iii.) If $\psi: \mathcal{U} \rightarrow \mathcal{U}$ is a local chart with $\psi(p) = (x^1(p), \dots, x^n(p))$ then the tangent vectors (X_1, \dots, X_n) defined for any $f \in C^\infty(M; \mathbb{R})$ through the equation:

$$X_\mu [f] \Big|_p = \frac{\partial (f \circ \psi^{-1})}{\partial x^\mu} \Big|_p$$

form a basis - the coordinate basis - in $T_p M$.

Outline of today's lecture

2.3 Vector fields - the tangent bundle

2.4 Covector fields - the cotangent bundle

2.5 Tensors

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From now on, we will use a simplifying notation:

$$\bullet X_{\mu}|_p [f] = \frac{\partial (f \circ \varphi^{-1})}{\partial x^{\mu}} \Big|_{\varphi(p)} = \frac{\partial f}{\partial x^{\mu}} \Big|_p$$

$$\bullet V_p [x^{\mu}] \equiv V_p [x^{\mu} \circ \varphi] = V_p^{\mu}$$

• Einstein summation convention

$$\sum_{\mu=1}^n V_p [x^{\mu}] X_{\mu}|_p = V_p^{\mu} \frac{\partial}{\partial x^{\mu}} \Big|_p$$

2.3 Vector fields - the tangent bundle

tangent bundle: We call the union

$TM := \bigcup_{p \in M} \{p\} \times T_p M$ the tangent bundle of M .

vector fields We call a map

$V: M \rightarrow TM; p \mapsto T_p M$ a smooth vector field (symbolically denoted by $V \in \Gamma(TM)$) if

$V[f]: M \rightarrow \mathbb{R}; p \mapsto V_p[f]$ is a smooth function for all $f \in C^\infty(M; \mathbb{R})$

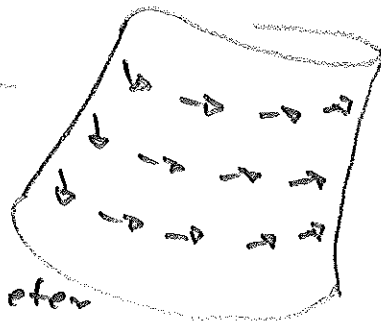
Now, the directional derivatives $X_p = \frac{\partial}{\partial x^k}$ are already smooth vectorfields; they form a basis hence:

V is a smooth vectorfield if (for a collection of charts covering all of M) its component functions $V^k = V[x^k]$ are smooth functions $V[x^k]: M \ni p \mapsto V_p[x^k] \in \mathbb{R}$ $\forall k=1, \dots, n$ themselves

The geometric interpretation of vectorfields

Vector fields define flows on the manifold

a flow defines a vectorfield



A flow is a smooth 1-parameter family $\phi_t : M \rightarrow M$ of diffeomorphisms:

- (i) ϕ_t is one-to-one $\forall t \in \mathbb{R}$
- (ii) $\phi_t \in \mathcal{C}^\infty(M; M)$
- (iii) $\phi_t^{-1} \in \mathcal{C}^\infty(M; M)$

If we pick a point $p \in M$; $t \mapsto \phi_t(p)$ defines a smooth curve, and we can then define a smooth vectorfield V through:

$$V_p = \left. \frac{d}{dt} \right|_{t=0} \phi_t(p) \in T_p M$$

i.e.

$$V_p[f] = \left. \frac{d}{dt} \right|_{t=0} (f \circ \phi_t)(p) \quad \forall f \in C^\infty(M; \mathbb{R})$$

the opposite is also true:
a vectorfield defines a flow

Given a vectorfield V we can always find its integral curves in M : We have to solve the following differential equation in every chart (U, ψ) covering M :

$$\frac{d}{dt} x^p(t) = V^p(x^1(t), \dots, x^n(t)) \quad \forall p=1, \dots, n$$

$$(x^1(0), \dots, x^n(0)) = p$$

and we then set:

$$\phi_t(p) = \psi^{-1}(x^1(t), \dots, x^n(t))$$

We can solve this differential equation iteratively, and thus write:

$$\phi_t = \exp(tV)$$

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Given two smooth vectorfields U, V we can define the Lie-Bracket:

$$\forall f \in C^\infty(M; \mathbb{R}):$$

$$[U, V]_p [f] = U_p [V [f]] - V_p [U [f]]$$

For any three smooth vectorfields we always have:

$$[[U, V], W] + [[U, W], V] + [[W, V], U] = 0$$

Given two commuting vectorfields $U, V: [U, V] = 0$ we can always find a two-parameter family of diffeomorphisms $\phi_{s,t}$ such that:

$$U_p = \left. \frac{d}{ds} \right|_{s,t=0} \phi_{s,t}(p) \in T_p M$$

$$V_p = \left. \frac{d}{dt} \right|_{s,t=0} \phi_{s,t}(p) \in T_p M$$

2.4 covector fields - the cotangent bundle

Given a (real) vector space V we can always build the dual vector space V^* as the space of all linear maps:

$$\omega : V \rightarrow \mathbb{R}; \quad V \mapsto \omega(V)$$

We have:

$$(i) \quad \forall a, b \in \mathbb{R}; \quad U, V \in V; \quad \omega \in V^*$$

$$\omega(aU + bV) = a\omega(U) + b\omega(V)$$

(ii) If $e_n, n=1, \dots, n$ is a basis in V then the equation

$$\tilde{e}^n(e_n) = \delta_n^n, \quad \tilde{e}^n \in V^*$$

defines a basis in V^* .

hence:

$$\dim(V) = \dim(V^*)$$

$$(iii) \quad (V^*)^* \cong V$$

At each point p of the manifold we have the tangent vector space $T_p M$. We now build its dual vector space $T_p^* M$ and call it the cotangent space $T_p^* M$.

We call the elements ω_p of $T_p^* M$ covariant vectors.

Co-vector bundle

- We call the union $T^* M = \bigcup_{p \in M} \{p\} \times T_p^* M$ the cotangent bundle of M .
- A map $\omega : M \rightarrow T^* M$ is a smooth covector field if for all smooth vector fields V the function $\omega(V) : M \rightarrow \mathbb{R}; p \mapsto \omega_p(V)$ is smooth. We then write $\omega \in \Gamma(T^* M)$.

Example of a smooth co-vector field

Let $f \in C^\infty(M; \mathbb{R})$ be a smooth function, we define its differential df_p at the point $p \in M$ by declaring for all tangent vectors $V_p \in T_p M$:

$$(df)_p(V) = V_p[f]$$

Now, for all $a, b \in \mathbb{R}$; $U, V \in T_p M$:

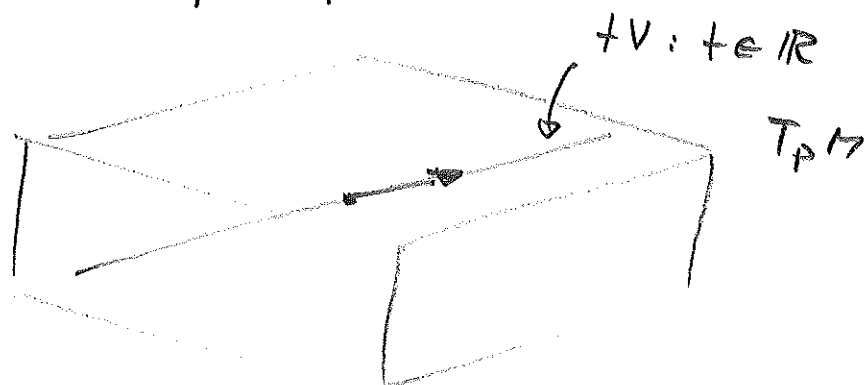
$$(df)_p(aU + bV) = a(df)_p(U) + b(df)_p(V)$$

Hence

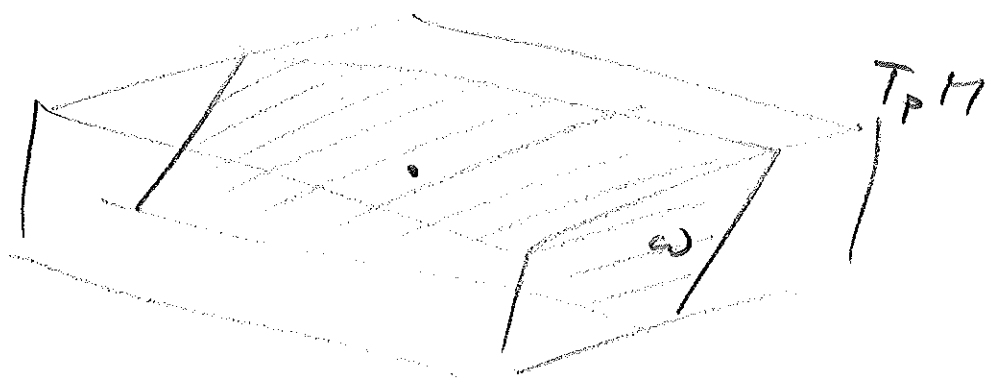
$$\boxed{df_p \in T_p^* M}$$

The geometric interpretation of co-vectors

- A vector $V \in T_p M$ defines a one-dimensional line in $T_p M$:



- A co-vector $\omega \in T_p^* M$ defines a $(n-1)$ -dimensional hypersurface $\{V \in T_p M \mid \omega(V) = 0\}$ in $T_p M$:



Note bene

- Vectorfields can always be integrated

$$V \rightarrow \exp(tV)$$

- For covectors this is not true:

$$\omega \in T^*M \text{ in general } \nexists f: \omega = df$$

A basis is T_p^*M

We have already seen that given a chart (\mathcal{U}, ψ) around $p \in \mathcal{U}$; $\psi(p) = (x^1(p), \dots, x^n(p))$ the derivatives:

$$(X_1|_p, \dots, X_n|_p) = \left(\frac{\partial}{\partial x^1} \Big|_p, \dots, \frac{\partial}{\partial x^n} \Big|_p \right)$$

form a basis in T_pM .

We now use them to construct a basis in T_p^*M :

Consider the differentials of the coordinate functions x^n $n=1, \dots, n$ around p :

$$(dx^n)_p(X_\nu) = X_\nu|_p[x^n] = \frac{\partial x^n}{\partial x^\nu} \Big|_p = \delta_\nu^n$$

We now set for any $\omega \in T_p^*M$:

$$\omega_\mu := \omega \left(\frac{\partial}{\partial x^\mu} \right)$$

and call ω_μ the μ -th component of ω w.r.t. the chart $(\mathcal{U}, \psi, \varphi)$ around p .

We now have:

$$\boxed{\omega = \omega_\mu dx^\mu}$$

Which can be seen as follows:

$$\omega(V) = \omega \left(V^\mu \frac{\partial}{\partial x^\mu} \right) = V^\mu \omega \left(\frac{\partial}{\partial x^\mu} \right) = V^\mu \omega_\mu$$

$$(\omega_\mu dx^\mu)(V) = \omega_\mu (dx^\mu)(V) = \omega_\mu V[x^\mu] = \omega_\mu V^\mu$$

Hence

$$\dim(T_p M) = \dim(T_p^* M) = n.$$

2.5 Tensor fields

- Given a (real) vector space V we saw how to construct its dual vector space V^* .
- We naturally have $(V^*)^* = V$.
- We can now also build tensors of rank (r, s) :

A tensor T of rank (r, s) is a multilinear map $T: \underbrace{V^* \times \dots \times V^*}_{r\text{-times}} \times \underbrace{V \times \dots \times V}_{s\text{-times}} \rightarrow \mathbb{R}$.

Basic properties

(i) The space of all tensors of rank (r,s) is a real vector space $V^{r,s}$:

$$\text{if } T, S \in V^{r,s}; a, b \in \mathbb{R}$$

$$aT + bS \in V^{r,s}$$

(ii) We naturally have

$$V^0_1 = V^*$$

$$V^1_0 = V$$

(iii) $\dim(V^{r,s}) = n^{r+s}$

Tensor product

If S is a tensor of rank (r, s) and T is a tensor of rank (l, m) we can build a tensor $(T \otimes S)$ of rank $(r+l, s+m)$ by declaring:

$$\begin{aligned}(S \otimes T)(\omega_1, \dots, \omega_{r+l}, v_1, \dots, v_{s+m}) &:= \\ &= S(\omega_1, \dots, \omega_r, v_1, \dots, v_s) T(\omega_{r+1}, \dots, \omega_{r+l}, v_{s+1}, \dots, v_{s+m})\end{aligned}$$

Basic properties

(i) $a, b \in \mathbb{R}$; R, S, T tensors:

- $(aR + bS) \otimes T = a(R \otimes T) + b(S \otimes T)$
- $T \otimes (aR + bS) = a(T \otimes R) + b(T \otimes S)$
- $a(T \otimes S) = (aT) \otimes S = T \otimes (aS)$

(ii) S, T tensors

$$S \otimes T \neq T \otimes S$$

(iii) R, S, T tensors

$$(R \otimes S) \otimes T = R \otimes (S \otimes T) = \\ = R \otimes S \otimes T$$

(iv) If $\{e_\mu\}$ $\mu=1, \dots, n$, $n = \dim(V)$ is a basis in V and $\{\tilde{e}^\mu\}$, $\mu=1, \dots, n$ is the dual basis in V^* then

$$e_{\mu_1} \otimes \dots \otimes e_{\mu_r} \otimes \tilde{e}^{\nu_1} \otimes \dots \otimes \tilde{e}^{\nu_s}$$

is a basis for all tensors of rank (r, s) , which follows immediately from:

$$(e_{\mu_1} \otimes \dots \otimes e_{\mu_r} \otimes \tilde{e}^{\nu_1} \otimes \dots \otimes \tilde{e}^{\nu_s})(\tilde{e}^{\kappa_1}, \dots, \tilde{e}^{\kappa_r}, e_{\beta_1}, \dots, e_{\beta_s}) = \\ = \delta_{\mu_1}^{\kappa_1} \dots \delta_{\mu_r}^{\kappa_r} \delta_{\beta_1}^{\nu_1} \dots \delta_{\beta_s}^{\nu_s}$$

We can thus write for any tensor of rank (n, s) :

$$T = T^{\mu_1 \dots \mu_n \nu_1 \dots \nu_s} e_{\mu_1} \otimes \dots \otimes e_{\mu_n} \otimes e^{\nu_1} \otimes \dots \otimes e^{\nu_s}$$

And can now define the following operations:

• Symmetrisation:

$$T_{[\mu_1 \dots \mu_k]} = \frac{1}{k!} \sum_{\pi \in S_k} T_{\mu_{\pi(1)} \dots \mu_{\pi(k)}}$$

• Anti-Symmetrisation:

$$T_{[\mu_1 \dots \mu_k]} = \frac{1}{k!} \sum_{\pi \in S_k} \text{sign}(\pi) T_{\mu_{\pi(1)} \dots \mu_{\pi(k)}}$$

• Contraction

If T is a tensor of rank (r, s) we define the contraction C^i_j over the i -th and j -th indices as the following operation:

$$C^i_j T = T^{n_1 \dots \alpha \dots n_n}_{n_1 \dots \alpha \dots n_s} \otimes e_{n_1} \otimes \dots \otimes e_{n_{i-1}} \otimes e_{n_{i+1}} \otimes \dots \otimes e_{n_n} \otimes \tilde{e}^{n_1} \otimes \dots \otimes \tilde{e}^{n_{j-1}} \otimes \tilde{e}^{n_{j+1}} \otimes \dots \otimes \tilde{e}^{n_s}$$

- $C^i_j T$ is a tensor of rank $(n-1, s-1)$

The result of the operations: contraction, symmetrization and anti-symmetrization does not depend on the actual basis chosen.